

Zero-resistance superconducting phase in BaFe₂As₂ under high pressureFumihiro Ishikawa,^{1,*} Naoya Eguchi,² Michihiro Kodama,² Koji Fujimaki,¹ Mari Einaga,¹ Ayako Ohmura,³ Atsuko Nakayama,³ Akihiro Mitsuda,⁴ and Yuh Yamada²¹Graduate School of Science and Technology, Niigata University, Niigata 950-2181, Japan²Department of Physics, Niigata University, Niigata 950-2181, Japan³Center for Transdisciplinary Research, Niigata University, Niigata 950-2181, Japan⁴Department of Physics, Kyushu University, Fukuoka 812-8581, Japan

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A single-crystalline sample of BaFe₂As₂ was prepared by high-pressure synthesis, and temperature dependence of resistivity was measured under high pressures of up to 9.0 GPa. Application of 3.0 GPa pressure suppresses the spin-density wave transition and induces superconductivity with the zero resistance state at a superconducting transition temperature $T_c=35$ K. T_c decreases monotonically with increasing pressure. The appearance of the superconductivity and the disappearance of the anomaly due to the spin-density wave transition occur simultaneously at a pressure of 3.0 GPa. At 3.0 GPa, a superconducting transition is induced in BaFe₂As₂; however, the superconductivity is suppressed at higher pressure.

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Since the discovery of superconductivity in LaFeAsO_{1-x}F_x at 26 K,¹ various superconductors with an iron-arsenide layer have been reported. Potassium-doped Ba_{0.55}K_{0.45}Fe₂As₂ with a ThCr₂Si₂-type structure, which also has an iron-arsenide layer, exhibits a superconducting transition at 38 K.² However, nondoped BaFe₂As₂ does not exhibit superconductivity but does have a spin-density wave (SDW) transition at 140 K accompanied by a structural transformation from tetragonal (*I4/mmm*) at higher temperature to an orthorhombic (*Fmmm*) structure at lower temperature.^{3,4}

For the BaFe₂As₂ system, the substitution of K for Ba suppresses the SDW and structural transitions and induces superconductivity at low temperatures.^{2,3} For nondoped BaFe₂As₂, application of pressure also suppresses both transitions, and superconductivity appears at low temperatures. The Meissner effect is observed in BaFe₂As₂ under pressure by measurements of dc magnetization between 2.5 and 6.0 GPa.⁵ The pressure dependence of the superconducting transition temperature T_c has a maximum value of 29 K near 4 GPa. The resistivity of BaFe₂As₂ drops below 20–30 K under high pressures above 3 GPa.⁶ However, no zero resistance appeared under pressures of up to 13 GPa, while zero resistivity under high pressure was reported by various groups for the AFe₂As₂ (A=Ca,Sr) systems.⁷⁻¹⁰

Zero resistivity observations are important to clarify the intrinsic properties of pressure-induced superconductivity in the BaFe₂As₂ system. In this Brief Report, we present recent results of high-pressure synthesis and observation of zero resistance for the pressure-induced superconductivity in BaFe₂As₂.

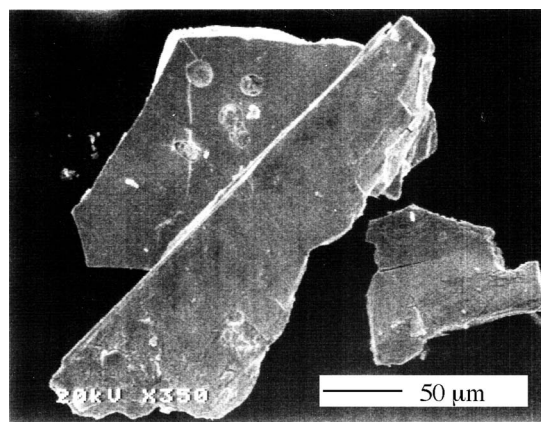
Samples were prepared by the high-pressure synthesis method using a cubic-anvil-type apparatus (TRY-00120, Try Engineering Co. Ltd). Pure Ba(>99%) and FeAs (99.5%) starting materials were mixed in appropriate amounts to prepare the nominal BaFe₂As₂ composition. Synthesis of BaFe₂As₂ was achieved by heating the mixture in a BN crucible under a pressure of approximately 3 GPa at 1473 K for 2 h. A scanning electron microscope (SEM) image of the resulting sample is shown in Fig. 1. Powder x-ray diffraction (XRD) patterns measured using Cu $K\alpha$ radiation confirmed

that the sample is single phase with the (001) orientation on the surface, which indicates a single-crystalline sample.

Pressures of up to 9.0 GPa were produced using a modified Bridgman anvil cell with a Teflon capsule.^{11,12} The Teflon capsule has a sample space with an inner diameter of 1.5 mm. A mixture of Fluorinert FC-70 and 77 was used as the pressure-transmitting medium. Pressure was estimated from the load-pressure calibration curve measured in advance.¹² The calibration curve was obtained from the resistive transition of Bi due to the *I-II*, *II-III*, and *III-V* transitions at room temperature and the superconductive transition temperature of Pb measured by ac susceptibility at low temperature. Difference in the load-pressure curves between room and low temperatures is less than 0.1 GPa at same load.

The modified Bridgman-anvil-type cell was cooled using a liquid He-free-type 4 K cryocooler (Iwatani Industrial Gases Corp.). The electrical resistance under high pressures was measured by a conventional four-probe method. As electrical leads, gold wires, 20 μ m in diameter, were used with a silver-loaded epoxy resin (Eccobond 56c, Emerson and Cuming) on the specimen.

The temperature dependence of resistivity in BaFe₂As₂ under high pressures up to 9.0 GPa is shown in Fig. 2. The

FIG. 1. SEM image of the BaFe₂As₂ sample.

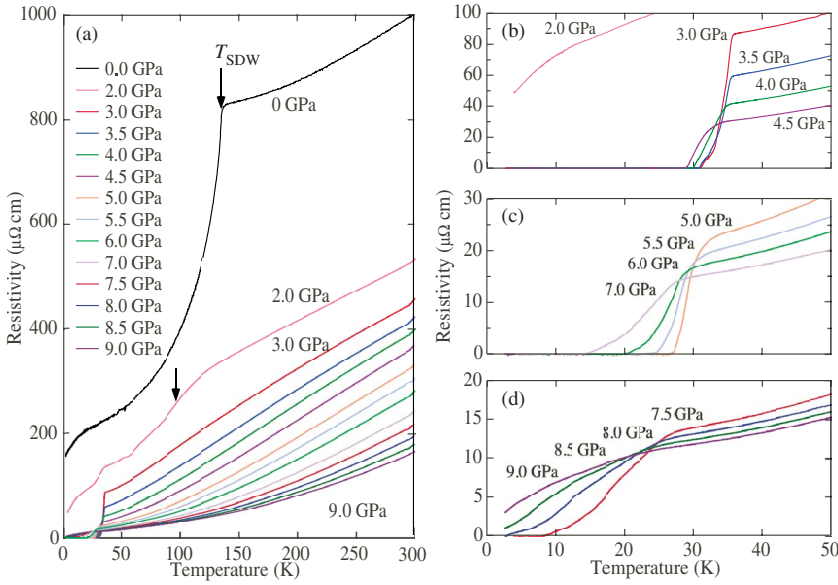


FIG. 2. (Color online) Temperature T dependence of the resistivity ρ in BaFe_2As_2 under pressure P ; (a) below 300 K and at $0 \leq P \leq 9.0$ GPa, and below 50 K at (b) $2.0 \leq P \leq 4.5$ GPa, (c) $5.0 \leq P \leq 7.0$ GPa, and (d) $7.5 \leq P \leq 9.0$ GPa. Arrows indicate the SDW transition temperature derived from the peak of the temperature derivative of resistivity, $d\rho/dT$.

superconducting transition, with zero resistance, appears at 3.0 GPa, as shown in Figs. 2(a) and 2(b). At ambient pressure, an anomaly due to the SDW transition³ appears at around $T_{\text{SDW}}=134$ K, which is derived from the peak of the temperature derivative of resistivity, $d\rho/dT$ (see the inset in Fig. 3). This temperature agrees well with the value of 131 K reported by Fukazawa *et al.*⁶ At increased pressure of 2.0 GPa, the anomaly becomes broader and T_{SDW} decreases to 96 K. At 3.0 GPa, the superconducting transition with zero resistance occurs at $T_c=35$ K, which is defined as the temperature of the onset of superconductivity. The superconducting transition is very sharp between 3.0 and 5.0 GPa, as shown in Figs. 2(b) and 2(c); the difference between T_c and the temperature where zero resistance state appears $T_{c-\text{zero}}$, is less than 5 K at 3.0 GPa. With increasing pressure, the transition temperature decreases monotonically, and the differ-

ence between T_c and $T_{c-\text{zero}}$ increases. At pressures above 7.5 GPa, no clear transition is observed and the zero resistance state disappears, as shown in Fig. 2(d). In more increased pressures of up to 9.0 GPa, $T_{c-\text{zero}}$ is not observed even at the lowest temperature in the present study, which is about 3 K. These results suggest that a pressure of 3.0 GPa induces the superconducting transition in BaFe_2As_2 , whereas higher pressures simply suppress superconductivity. It is noted that the appearance of the superconductivity and the disappearance of the anomaly due to the SDW transition occur simultaneously at 3.0 GPa, as shown in Fig. 2(a). At 2.0 GPa, the SDW anomaly at 96 K and the resistivity drop at 35 K are observed; however, no zero resistance state appears, even at 3 K. Thus, the suppression of the SDW state is necessary for the appearance of the superconducting state with zero resistance in BaFe_2As_2 .

Magnetization measurements of BaFe_2As_2 under high pressures indicate that the superconducting transition occurs gradually with increasing pressure, while the superconducting transition in SrFe_2As_2 occurs abruptly, accompanied by a maximum in superconducting temperature.⁵ The present results for the pressure dependence of T_c for BaFe_2As_2 are similar to those for SrFe_2As_2 (Ref. 5) and CaFe_2As_2 (Ref. 8) in terms of the appearance of superconductivity; the onset of superconductivity occurs abruptly at a critical pressure.

The pressure dependences of T_c , $T_{c-\text{zero}}$, and T_{SDW} are presented in Fig. 3. T_c and $T_{c-\text{zero}}$ decrease with increasing pressure, and both the SDW anomaly and superconductivity are suppressed by pressure. The pressure dependence of T_c shown in Fig. 3 is very similar to that for the characteristic temperature of polycrystalline BaFe_2As_2 .⁶ However, for the polycrystalline sample, the resistivity drops sharply at the characteristic temperature, which probably corresponds to T_c , at the pressure of 13.0 GPa; the difference between T_c and $T_{c-\text{zero}}$ is small even in the pressures above 9 GPa. Moreover, the pressure dependence of T_c obtained from the magnetization measurements up to the pressure of 6 GPa (Ref. 5) shows little change in T_c between 5 and 6 GPa. Thus, T_c and $T_{c-\text{zero}}$ behavior shown in Fig. 3 is not consistent with the previous reports,^{5,6} especially in high-pressure region above 5 GPa.

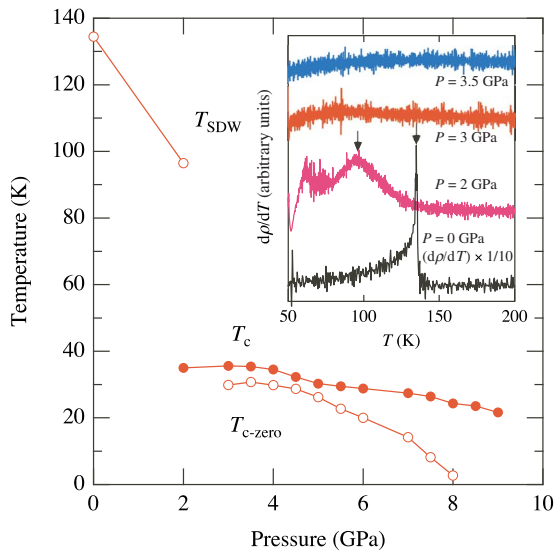


FIG. 3. (Color online) T - P phase diagram of BaFe_2As_2 determined from resistivity anomalies under high pressure. Inset shows $d\rho/dT$ vs T in BaFe_2As_2 under high pressure. Arrows indicate the SDW anomaly derived from the peak of $d\rho/dT$.

Several authors reported the effect of pressure on the superconductivity of the carrier-doped $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ and $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ systems.^{13,14} Resistivity measurements of $\text{BaFe}_{1.92}\text{Co}_{0.08}\text{As}_2$ under high pressures show that the SDW transition is not completely suppressed up to 2.5 GPa and T_c increases slightly with increasing pressure.¹³ Moreover, T_c is almost independent of pressure in $\text{BaFe}_{1.8}\text{Co}_{0.2}\text{As}_2$, which shows no anomaly due to the SDW transition.¹³ In contrast, an opposite effect of pressure on T_c was reported for the $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ system; T_c decreases with increasing pressure.¹⁴ Resistivity measurements of $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ showed that the SDW state is suppressed by K doping. Thus, in the BaFe_2As_2 and related systems, T_c decreases with increasing pressure when the SDW transition is completely suppressed by pressure or doping. Consequently, T_c on these systems under high pressure depends on the competition between two kinds of pressure effects on superconductivity: direct suppression effect and enhancement effect via suppression of the SDW state. These results may support such a scenario as follows: (1) the superconductivity of BaFe_2As_2 and related systems is suppressed by the SDW state. (2) The suppression of the SDW state by application of pressure re-

sults in onset of the superconductivity. (3) Further increase in pressure suppresses the superconductivity after the SDW state is completely suppressed.

In summary, a single-crystalline sample of BaFe_2As_2 was prepared by high-pressure synthesis using a cubic-anvil-type apparatus, and the temperature dependence of resistivity was measured under high pressures up to 9.0 GPa. Application of 3.0 GPa pressure suppresses the SDW transition and induces superconductivity with the zero resistance state at $T_c = 35$ K. T_c decreases monotonically with increasing pressure, and no clear transition is observed above 7.5 GPa. The appearance of the superconductivity and the disappearance of the SDW transition anomaly occur simultaneously at a pressure of 3.0 GPa. The effect of pressure suppresses both the SDW state and superconductivity in BaFe_2As_2 .

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*fumihito@phys.sc.niigata-u.ac.jp

¹Y. Kamihara, T. Watanabe, M. Hirano, and H. Hosono, *J. Am. Chem. Soc.* **130**, 3296 (2008).

²M. Rotter, M. Tegel, and D. Johrendt, *Phys. Rev. Lett.* **101**, 107006 (2008).

³M. Rotter, M. Tegel, D. Johrendt, I. Schellenberg, W. Hermes, and R. Pöttgen, *Phys. Rev. B* **78**, 020503(R) (2008).

⁴Q. Huang, Y. Qiu, W. Bao, M. A. Green, J. W. Lynn, Y. C. Gasparovic, T. Wu, G. Wu, and X. H. Chen, *Phys. Rev. Lett.* **101**, 257003 (2008).

⁵P. L. Alireza, Y. T. C. Ko, J. Gillett, C. M. Petrone, J. M. Cole, G. G. Lonzarich, and S. E. Sebastian, *J. Phys.: Condens. Matter* **21**, 012208 (2009).

⁶H. Fukazawa, N. Takeshita, T. Yamazaki, K. Kondo, K. Hirayama, Y. Kohori, K. Miyazawa, H. Kito, H. Eisaki, and A. Iyo, *J. Phys. Soc. Jpn.* **77**, 105004 (2008).

⁷H. Kotegawa, H. Sugawara, and H. Tou, *J. Phys. Soc. Jpn.* **78**, 013709 (2008).

⁸M. S. Torikachvili, S. L. Bud'ko, N. Ni, and P. C. Canfield, *Phys. Rev. Lett.* **101**, 057006 (2008).

⁹T. Park, E. Park, H. Lee, T. Klimczuk, E. D. Bauer, F. Ronning, and J. D. Thompson, *J. Phys.: Condens. Matter* **20**, 322204 (2008).

¹⁰K. Igawa, H. Okada, H. Takahashi, S. Matsuishi, Y. Kamihara, M. Hirano, H. Hosono, K. Matsubayashi, and Y. Uwatoko, *J. Phys. Soc. Jpn.* **78**, 025001 (2009).

¹¹T. Nakanishi, N. Takeshita, and N. Mōri, *Rev. Sci. Instrum.* **73**, 1828 (2002).

¹²F. Ishikawa, K. Fukuda, S. Sekiya, A. Kaeriyama, Y. Yamada, and A. Matsushita, *J. Phys. Soc. Jpn.* **76**, Suppl. A, 92 (2007).

¹³K. Ahilan, J. Balasubramanian, F. L. Ning, T. Imai, A. S. Sefat, R. Jin, M. A. McGuire, B. C. Sales, and D. Mandrus, *J. Phys.: Condens. Matter* **20**, 472201 (2008).

¹⁴M. S. Torikachvili, S. L. Bud'ko, N. Ni, and P. C. Canfield, *Phys. Rev. B* **78**, 104527 (2008).